

Law of universal gravitation with finite velocity of gravity and mathematical model of motion of a finite number of material points

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Abstract – The law of universal gravitation is introduced taking into account the finiteness of the gravitational velocity. Based on this law, a mathematical model of the motion of a finite number of material points is constructed, a separate case of which is the classical model of the motion of points, which is described by a system of ordinary differential equations. The constructed model is a system of nonlinear differential equations with deviating argument and functional equations. It more accurately describes the dynamics of a finite number of material points than the corresponding classical model. A mathematical model of the motion of two material points is also considered.

Keywords: law of universal gravitation with finite velocity of gravity, mathematical model of motion of a finite number of material points.

I. INTRODUCTION

In Newton's theory of gravity, the gravity velocity c_g is not included in any formula and is considered to be infinitely large ($c_g = \infty$).

In A. Einstein's general theory of relativity for the speed of gravity, it is assumed that $c_g = c$, where c is the speed of light [1], [2].

The speed of gravity can be estimated by the transmission rate of the influence of the gravitational field on the measurement results. So, in 2002, Kopeikin and Fomalont conducted an experiment based on radio interferometry with an extremely long base, in which radiation from a distant quasar QSO J0842+1835, which passed near a massive body - Jupiter, was recorded by Earth's radio telescopes [3]. Analysis of the experimental data yielded a speed of gravity close in magnitude to the speed of light, with an accuracy of about 20 %.

The second method of measuring the speed of gravity is associated with the fixation of gravitational waves from distant stellar sources simultaneously with a light signal.

The first such measurement was obtained in 2017 for the GW170817 gravitational wave generated by the fusion of two neutron stars (the distance to the source was 13×10^7 light years) using laser-interferometric gravitational-wave detectors of the LIGO-Virgo detector grid. Behind this event, the deviation of the speed of gravitational waves from the speed of light, if such a deviation exists, ranges from -3×10^{-15} to $+0.7 \times 10^{-15}$, those are compatible with zero within the error [4].

This article is devoted to the law of universal gravitation, taking into account the finiteness of the speed of gravity and its application to the construction of a mathematical model of motion of a finite number of bodies. Such a model is not a system of ordinary differential equations, as in classical celestial mechanics [5], but a system of differential equations with a deviating argument and functional equations, as in [6]. The dynamics of bodies, taking into account the finiteness of the gravity velocity, is characterized by properties that are impossible in the case of $c_g = \infty$.

II. THE LAW OF GRAVITY WITH FINITE SPEED OF GRAVITY

Consider two material points M_1 and M_2 with masses m_1 and m_2 respectively. According to the law of universal gravitation and Newton's second law, these points move in space. The movement of the points will be considered with respect to the rectangular coordinate system x, y, z with the origin at some point O . We consider that the coordinate system is inertial and only gravity generated by another point acts on each point. The position of the points M_1 and M_2 at time t is determined by their vector functions $\vec{r}_1(t)$ and $\vec{r}_2(t)$.

To study the motion of points M_1 and M_2 , it is necessary to know the forces with which each of these points attracts the other.

If the speed of gravity were infinite, as in Newton's theory, then, on the basis of the law of universal gravitation,

at the moment t , the point M_2 would attract the point M_1 with force

$$\vec{F}_{2,1,\infty}(t) = \frac{Gm_1m_2}{|\vec{r}_2(t) - \vec{r}_1(t)|^3} (\vec{r}_2(t) - \vec{r}_1(t)), \quad (1)$$

where G is the gravitational constant and $|\vec{r}_2(t) - \vec{r}_1(t)|$ is the Euclidean length of the vector $\vec{r}_2(t) - \vec{r}_1(t)$. The direction of this force coincides with the direction of vector $\vec{r}_2(t) - \vec{r}_1(t)$.

Similarly, at time t , point M_1 attracted point M_2 with force

$$\vec{F}_{1,2,\infty}(t) = \frac{Gm_1m_2}{|\vec{r}_1(t) - \vec{r}_2(t)|^3} (\vec{r}_1(t) - \vec{r}_2(t)).$$

According to paragraph 1, the speed of gravity is finite. Therefore, the action of one point on another is carried out taking into account the delay of the gravitational field.

Let us explain the effect of the delay of the gravitational field on the forces with which the considered points 1 and 2 are attracted.

Based on the finiteness of the speed of gravity, a another force acts on point M_1

$$\begin{aligned} \vec{F}_{2,1,c}(t) &= \\ &= \frac{Gm_1m_2}{|\vec{r}_2(t - \tau_{2,1}(t)) - \vec{r}_1(t)|^3} (\vec{r}_2(t - \tau_{2,1}(t)) - \vec{r}_1(t)), \quad (2) \end{aligned}$$

where the delay of gravity $\tau_{2,1}(t)$ in (2) satisfies the ratio of

$$c\tau_{2,1}(t) = |\vec{r}_2(t - \tau_{2,1}(t)) - \vec{r}_1(t)| \quad (3)$$

and c the speed of gravity.

Indeed, let points M_2 and M_1 move along curves, parts of which are shown in Fig. 1, with speeds

$$\vec{v}_2(t) = \frac{d\vec{r}_2(t)}{dt} \quad \text{and} \quad \vec{v}_1(t) = \frac{d\vec{r}_1(t)}{dt}.$$

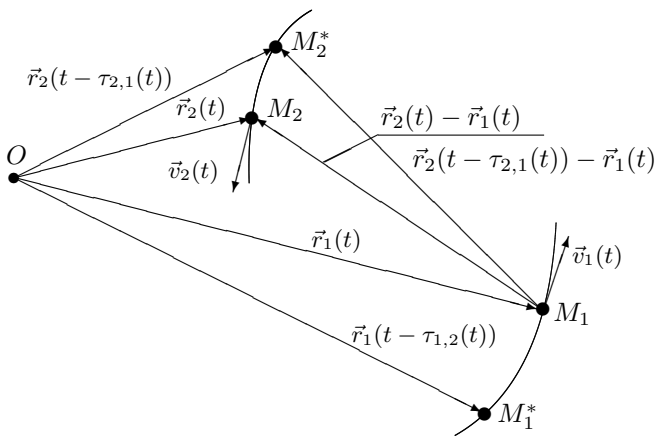


Fig. 1. Points M_1 and M_2 at times t and $t - \tau_{2,1}(t)$.

Let point M_2 at time $t - \tau_{2,1}(t)$, where $\tau_{2,1}(t)$ satisfies

(3), coincide with point M_2^* . Time interval $[t - \tau_{2,1}(t), t]$ corresponds to curve $M_2^*M_2$, along which point M_2 moves. This time interval is sufficient for the gravitational field with a velocity of c to propagate from the point M_2^* to the point M_1 . Thus, at time t , point M_1 is not affected by force (1), but by force (2).

Note that the attracting point for point M_1 at time t is not point M_2 , but point M_2^* .

Similarly, based on the finiteness of the speed of gravity, force

$$\begin{aligned} \vec{F}_{1,2,c}(t) &= \\ &= \frac{Gm_1m_2}{|\vec{r}_1(t - \tau_{1,2}(t)) - \vec{r}_2(t)|^3} (\vec{r}_1(t - \tau_{1,2}(t)) - \vec{r}_2(t)) \quad (4) \end{aligned}$$

acts on point M_2 and the attracting point for point M_2 at time t is not point M_1 , but point M_1^* , which is determined by vector $\vec{r}_1(t - \tau_{1,2}(t))$.

The delay of gravity $\tau_{1,2}(t)$ in equality (4) satisfies relation

$$c\tau_{1,2}(t) = |\vec{r}_1(t - \tau_{1,2}(t)) - \vec{r}_2(t)|. \quad (5)$$

The existence of functions $\tau_{2,1}(t)$ and $\tau_{1,2}(t)$, which satisfy relations (3) and (5), is shown in [6]. Based on the implicit function theorem [7, p. 449–453], these functions are continuous and differentiable.

Note that the forces $\vec{F}_{2,1,c}(t)$ and $\vec{F}_{1,2,c}(t)$ may differ in magnitude and not be collinear. Also based on (3) and (5) for each point in time t

$$\lim_{c \rightarrow +\infty} \vec{F}_{2,1,c}(t) = \vec{F}_{2,1,\infty}(t)$$

and

$$\lim_{c \rightarrow +\infty} \vec{F}_{1,2,c}(t) = \vec{F}_{1,2,\infty}(t).$$

Thus, Newton's law of universal gravitation opens the way to a more general law (see (2)–(5)) and coincides with it in the extreme case (with $c = +\infty$).

Note that to use formulas (2) and (4), information on vector functions $\vec{r}_1(t)$ and $\vec{r}_2(t)$ and scalar functions $\tau_{2,1}(t)$ and $\tau_{1,2}(t)$ are needed. These functions can be found using differential equations with a deviating argument, which describe the motion of points M_1 and M_2 , and are obtained using Newton's second law (see [8]).

III. THE MATHEMATICAL MODEL OF THE MOTION OF A FINITE NUMBER OF MATERIAL POINTS

We pay attention to the mathematical model of the movement of points M_i , $i = \overline{1, n}$, with masses m_i , $i = \overline{1, n}$, respectively, where $n \geq 2$. The movement of these points will be considered relative to the rectangular coordinate system x, y, z with the origin at some point O . We assume that this system is inertial.

The movement of points M_i , $i = \overline{1, n}$, can be described by some vector functions $\vec{r}_i(t)$, $i = \overline{1, n}$, respectively.

Fix arbitrary points M_i and M_j , where $i, j \in \{1, 2, \dots, n\}$ and $i \neq j$. In accordance with paragraph 2, point M_j attracts point M_i with force

$$\begin{aligned} & \vec{F}_{j,i,c}(t) = \\ & = \frac{Gm_j m_i}{|\vec{r}_j(t - \tau_{j,i}(t)) - \vec{r}_i(t)|^3} (\vec{r}_j(t - \tau_{j,i}(t)) - \vec{r}_i(t)), \end{aligned} \quad (6)$$

where the delay of gravity $\tau_{j,i}(t)$ in equality (6) satisfies relation

$$c\tau_{j,i}(t) = |\vec{r}_j(t - \tau_{j,i}(t)) - \vec{r}_i(t)|. \quad (7)$$

We use force $\sum_{j \in N_{i,n}} \vec{F}_{j,i,c}(t)$, which acts on point M_i generated by points M_j , $j \in N_{i,n}$, where

$$N_{i,n} = \{1, 2, \dots, n\} \setminus \{i\}, \quad i = \overline{1, n}.$$

Based on Newton's second law for vector functions $\vec{r}_i(t)$, $i = \overline{1, n}$, the relations

$$m_i \frac{d^2 \vec{r}_i(t)}{dt^2} = \sum_{j \in N_{i,n}} \vec{F}_{j,i,c}(t), \quad i = \overline{1, n}.$$

are satisfied. It follows from relations (6) and (7) that vector functions $\vec{r}_i(t)$, $i = \overline{1, n}$, are solutions of the following system of differential equations with deviating argument and functional equations

$$\begin{cases} \frac{d^2 \vec{r}_i(t)}{dt^2} = \\ = \sum_{j \in N_{i,n}} \frac{Gm_j}{|\vec{r}_j(t - \tau_{j,i}(t)) - \vec{r}_i(t)|^3} (\vec{r}_j(t - \tau_{j,i}(t)) - \vec{r}_i(t)), & i = \overline{1, n}, \\ c\tau_{j,i}(t) = |\vec{r}_j(t - \tau_{j,i}(t)) - \vec{r}_i(t)|, & i = \overline{1, n}, j = \overline{1, n}, i \neq j. \end{cases} \quad (8)$$

Obviously, for each t the correct relations are

$$\lim_{c \rightarrow +\infty} \tau_{j,i}(t) = 0, \quad i = \overline{1, n}, j = \overline{1, n}, i \neq j,$$

$$\lim_{c \rightarrow +\infty} \vec{F}_{j,i,c}(t) = \vec{F}_{j,i,\infty}(t), \quad i = \overline{1, n}, j = \overline{1, n}, i \neq j,$$

where

$$\vec{F}_{j,i,\infty}(t) = \frac{Gm_j m_i}{|\vec{r}_j(t) - \vec{r}_i(t)|^3} (\vec{r}_j(t) - \vec{r}_i(t)).$$

Therefore, system (8) is a generalization of the classical model of the movement of points M_i , $i = \overline{1, n}$,

$$\begin{cases} \frac{d^2 \vec{r}_i(t)}{dt^2} = \sum_{j \in N_{i,n}} \frac{Gm_j}{|\vec{r}_j(t) - \vec{r}_i(t)|^3} (\vec{r}_j(t) - \vec{r}_i(t)), \\ i = \overline{1, n}, \end{cases} \quad (9)$$

which is obtained from (8) for $c = \infty$.

When finding the trajectories of movement of points M_i , $i = \overline{1, n}$, it is necessary to use, in addition to the equations of system (8), also initial or boundary conditions (see [6], [8], [9]). System (8), together with these conditions, is a mathematical model of the movement of points M_i , $i = \overline{1, n}$, that takes into account the finiteness of the speed of gravity.

Obviously, system (8), taking into account the initial conditions, can be used as a mathematical model of a stellar system with $n - 1$ planets with masses m_i , $i = \overline{2, n}$, that move in a force field generated by a star with mass m_1 and planets.

Note that system (8) contains n vector equations and $n(n - 1)$ scalar equations. Therefore, this system is equivalent to a system of scalar equations, the number of which coincides with

$$3n + n(n - 1) = n^2 + 2n = (n + 1)^2 - 1.$$

The classical model (9) of the motion of n material points, which contains n vector equations, is equivalent to a system of $3n$ scalar equations.

From this and the fact that

$$\lim_{n \rightarrow \infty} \frac{(n + 1)^2 - 1}{3n} = +\infty,$$

it follows that even if the complexity of the equations of system (8) is not taken into account, studying system (8) at large n is a more difficult task than studying system (9).

Despite this, it is possible to find such properties of system (8) that are not characteristic of system (9) (see [6], [9]).

IV. THE MATHEMATICAL MODEL OF THE MOVEMENT OF TWO POINTS

In case $n = 2$, the system of equations (8) takes the form

$$\begin{cases} \frac{d^2 \vec{r}_1(t)}{dt^2} = \\ = \frac{Gm_2}{|\vec{r}_2(t - \tau_{2,1}(t)) - \vec{r}_1(t)|^3} (\vec{r}_2(t - \tau_{2,1}(t)) - \vec{r}_1(t)), \\ \frac{d^2 \vec{r}_2(t)}{dt^2} = \\ = \frac{Gm_1}{|\vec{r}_1(t - \tau_{1,2}(t)) - \vec{r}_2(t)|^3} (\vec{r}_1(t - \tau_{1,2}(t)) - \vec{r}_2(t)), \\ c\tau_{2,1}(t) = |\vec{r}_2(t - \tau_{2,1}(t)) - \vec{r}_1(t)|, \\ c\tau_{1,2}(t) = |\vec{r}_1(t - \tau_{1,2}(t)) - \vec{r}_2(t)|. \end{cases} \quad (10)$$

This system, together with the initial or boundary conditions considered in [9], [11] is a mathematical model of

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the motion of two points with masses m_1 and m_2 . System (10) has been thoroughly investigated in the works [8–12].

V. ADDITIONAL REMARKS AND LITERATURE

The general case of a system of n bodies, taking into account the finite speed of gravity, was considered by the author in article [6]. This article also pays attention to the construction of a mathematical model of the solar system with a finite speed of gravity and to obtain some new properties for it.

In [9], the problem of two bodies of arbitrary masses was studied taking into account the speed of gravity. The non-Kaplerian behavior and instability of the motion of these bodies are also shown.

In [8], systems of differential equations with delays and constraints on delays and derivatives of solutions, which are used in [9] and [10], are studied.

Articles [13], [14], and [15] show the instability of unbounded trajectories of stellar systems.

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